

# Determination of exchange energies in the sawtooth spin ring $\{\text{Mo}_{75}\text{V}_{20}\}$ by ESR

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The magnetism of the polyoxometalate cluster  $\{\text{Mo}_{75}\text{V}_{20}\}$ , containing a sawtooth ring of 10 corner-sharing triangles located on the equator of the barrel-shaped molecule, has remained debatable since it is masked by contributions from impurities as well as temperature-independent paramagnetism. In this article we demonstrate the usefulness of ESR measurements since the temperature dependence of the ESR intensity can discriminate between impurity and molecular contributions. We determine the exchange parameters and therefore also the low-lying spectrum of  $\{\text{Mo}_{75}\text{V}_{20}\}$ , especially the low-lying singlet states which so far have been probed solely by specific heat measurements.

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## I. INTRODUCTION

The series of Keplerate molecules  $\{\text{Mo}_{72}\text{Fe}_{30}\}$ ,  $\{\text{Mo}_{72}\text{Cr}_{30}\}$ ,  $\{\text{Mo}_{72}\text{V}_{30}\}$  and  $\{\text{W}_{72}\text{V}_{30}\}$  is one of the beautiful creations made possible by recent developments of modern chemistry.<sup>1–10</sup> In these nanosized molybdate or tungstate-based molecules, 30 magnetic ions are located on the 30 vertices of an icosidodecahedron and are antiferromagnetically coupled, resulting in triangular and pentagonal networks. As a result, this leads to a strongly frustrated cluster with a huge number of quantum states.<sup>11</sup> Since the number of states,  $(2s+1)^{30}$ , can be varied by substituting the magnetic ions ( $2^{30}$  and  $6^{30}$  for V ( $s=1/2$ ) and Fe ( $s=5/2$ ), respectively), these are ideal systems for studying the transitions from quantum to classical behavior.

Among synthesized derivatives of the Keplerate clusters, the compound  $\{\text{Mo}_{75}\text{V}_{20}\}$  is analogous to  $\{\text{Mo}_{72}\text{V}_{30}\}$  except that the 10  $\text{V}^{4+}$  ions ( $s=1/2$ ) located adjacent to the north and south poles of the icosidodecahedron are substituted by nonmagnetic ions.<sup>12</sup> Therefore, this compound is equivalent to a sawtooth chain of 10 triangles with a periodic boundary condition as schematically shown in the inset of Fig. 1. Basically,  $\{\text{Mo}_{75}\text{V}_{20}\}$  is classified as a sawtooth spin ring rather than a Keplerate cluster since it has only a partial substructure of  $\{\text{Mo}_{72}\text{V}_{30}\}$ . However, determining the exchange couplings of  $\{\text{Mo}_{75}\text{V}_{20}\}$  can be a good reference for the characterization of its analogous compound  $\{\text{Mo}_{72}\text{V}_{30}\}$ . Moreover, sawtooth chains also belong to the class of frustrated antiferromagnetic spin systems with potentially very unusual magnetization curves,<sup>13,14</sup> as partly realized in the recently investigated magnetic material azurite.<sup>15,16</sup>

For  $\{\text{Mo}_{75}\text{V}_{20}\}$ , it turns out that the product,  $\chi T$ , of the magnetic susceptibility and temperature decreases steadily with decreasing temperature due to antiferro-

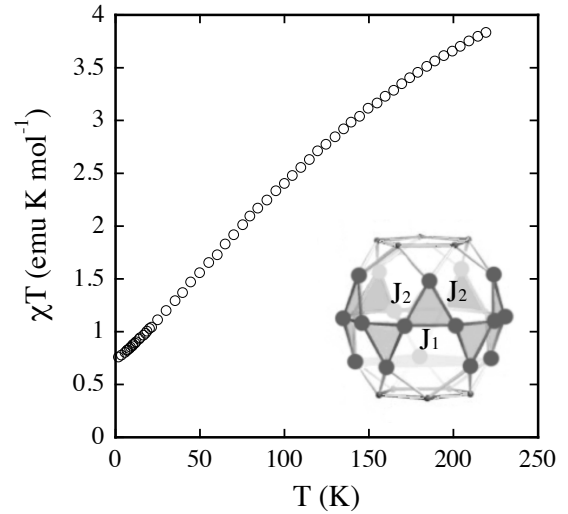


FIG. 1: Temperature dependence of  $\chi T$  from Ref. 3. The inset shows a scheme of the metal skeleton of  $\{\text{Mo}_{75}\text{V}_{20}\}$  highlighting the nearest-neighbor exchange pattern between the  $\text{V}^{4+}$  ions.

magnetic couplings, and it was assumed in Ref. 3 that the ground state is a singlet. Additionally, those authors estimated the exchange couplings between the magnetic ions from a theoretical fit to a reduced model that contains only a ring of six triangles. They concluded that their experimental results are well reproduced when  $J_1 = 288$  K,  $J_2 = 0.55 \cdot J_1$ , compare inset of Fig. 1, together with an additional coupling  $J_3 = 0.20 \cdot J_1$  between the tips of every second triangle.<sup>3</sup>

Takemura and Fukumoto refined those parameters by using a finite-temperature Lanczos method for the full system of 10 coupled triangles.<sup>17</sup> They found similar exchange coupling parameters, which are  $J_1 = 388$  K,

$J_2 = 0.42 \cdot J_1$ , and  $J_3 = 0.21 \cdot J_1$ . However, these values are still under discussion for various reasons. The main concern is the unrealistically large exchange  $J_3$  which should act across a distance of 11.7 Å, mediated by multi-center exchange pathways. Another problem is given by the fact that the susceptibility data are superimposed by an unknown amount of free vanadium ions (impurities) as well as by temperature-independent (i.e. van Vleck) paramagnetism typical of polyoxometalates.

We propose in this article that ESR is a unique method that overcomes these difficulties. It can separate the signals stemming from the intrinsic  $\{\text{Mo}_{75}\text{V}_{20}\}$  and from impurities since each ESR linewidth is qualitatively different. This enables us to obtain the pure magnetic response from  $\{\text{Mo}_{75}\text{V}_{20}\}$  in contrast to the susceptibility measurements in which the intrinsic and extrinsic responses are mixed. Since the temperature dependence of the ESR intensity is related to the energies of the excited states, we are able to obtain information about the exchange couplings. The ratio of  $J_1$  and  $J_2$  determines the frustration of the sawtooth chain of which the density of low-lying singlet states is a fingerprint.<sup>18,19</sup> Usually the density of singlet-levels is deduced from specific heat measurements, but we will show below that ESR is also capable of determining the density of low-lying singlet states.

The article is organized as follows. In Section II we briefly explain the experimental method. Section III contains our experimental results that are compared to the calculated ESR intensities assuming various sets of model parameters. The article closes with a short summary.

## II. EXPERIMENTAL

The Terahertz Electron Spin Resonance Apparatus in the Institute for Materials Research (TESRA-IMR) of Tohoku University has been used for the high field ESR measurements.<sup>20</sup> A simple transmission method with Faraday configuration has been employed. We used conventional Gunn oscillators for the millimeter wave radiation, and an InSb detector for transmission detection. In addition, a pulsed magnetic field up to 8.5 T can be generated from a 90 kJ capacitor bank. Two types of cryostats, a conventional  $^4\text{He}$  bath type cryostat and gas-flow type cryostat, were used for the measurements for the low and high temperature ranges, respectively. Powder samples were used in this study.

## III. RESULTS AND DISCUSSION

### A. High field ESR results

Figure 2 shows the typical temperature dependence of the ESR spectra for  $\{\text{Mo}_{75}\text{V}_{20}\}$ . The employed frequency is 190 GHz, and the temperature is varied from

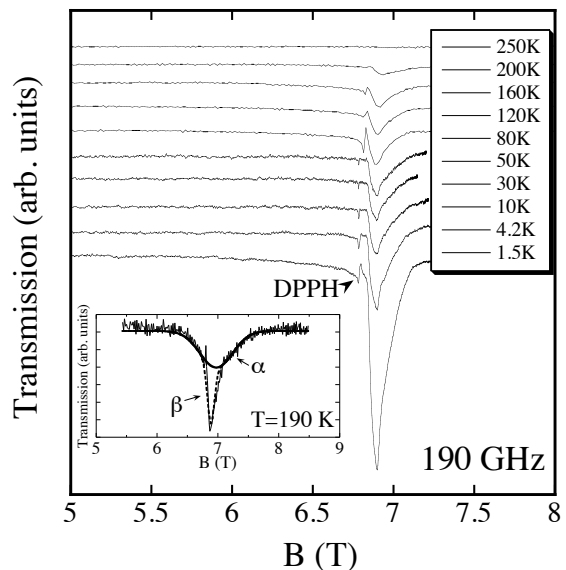


FIG. 2: Temperature dependence of ESR spectra for 190 GHz. The very narrow peak observed around 6.8 T is a field marker. The inset shows the spectrum for 190 K. The solid and broken lines are fitting curves for resonance  $\alpha$  and  $\beta$ , respectively (see text for details).

1.5 to 250 K. As shown in the inset of Fig. 2, two absorption lines, a broad and a sharp one, are clearly observed at 190 K. Hereafter, we denote the observed resonances as  $\alpha$  and  $\beta$  for the former and the latter, respectively. The tiny absorption observed at around 6.8 T is from DPPH, which stands for (2,2-diphenyl-1-picrylhydrazyl), a field marker. The intensity of  $\alpha$  gradually increases by decreasing the temperature, but quickly diminishes below 80 K, which is a typical ESR behavior of the excited states. On the other hand, the intensity of  $\beta$  is inversely proportional to the temperature, following Curie's law. Therefore,  $\alpha$  can be assigned to the resonance originating from  $\{\text{Mo}_{75}\text{V}_{20}\}$ , and  $\beta$  from the impurities at the cation sites.<sup>3</sup> Since the singlet ground state is ESR silent, the absence of intrinsic ESR signal from  $\{\text{Mo}_{75}\text{V}_{20}\}$  at low temperature and the ESR observation of the excited states for relatively high temperatures suggests that the ground state of  $\{\text{Mo}_{75}\text{V}_{20}\}$  is a singlet, which is consistent with magnetic susceptibility results.<sup>3</sup>

In the inset of Fig. 2 the absorption lines of  $\alpha$  and  $\beta$  are fitted by Gaussian curves (solid and broken curves, respectively), and integrated intensities for each temperature are obtained. The integrated intensity of impurities (i.e. of resonance  $\beta$ ),  $I_{\text{imp}}$ , versus temperature is shown as solid circles in Fig. 3. According to the Boltzmann distribution, the ESR intensity of the  $s = 1/2$  impurities can be written as

$$I_{\text{imp}} = w N_{\text{imp}} \tanh(\beta g \mu_B B / 2). \quad (1)$$

$w$  is a coefficient which is proportional to the power and frequency of the radiation,  $N_{\text{imp}}$  is the number

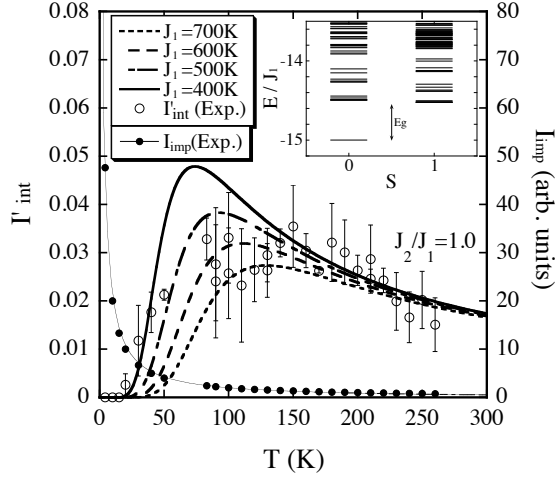


FIG. 3: Solid circles are the integrated intensities for resonance  $\beta$  coming from paramagnetic impurities; they are fitted using equation (1) (thin solid curve). Open circles are the normalized integrated intensities for resonance  $\alpha$  coming from the intrinsic  $\{\text{Mo}_{75}\text{V}_{20}\}$ . Thick solid and dashed curves are calculated curves of the normalized integrated intensity for various values of  $J_1$ . The calculated curves are obtained from the singlet and triplet energy levels (inset) assuming an equilateral triangle model. The energy is normalized by  $J_1$ .

of impurities,<sup>21</sup> and  $\beta = 1/k_B T$  the inverse temperature.  $g$  is the spectroscopic splitting factor and  $\mu_B$  the Bohr magneton, respectively. By using equation (1)  $wN_{\text{imp}} = 83.9 \pm 7.4$  is obtained from the fitting curve in Fig. 3 (thin solid curve).

On the other hand, the intrinsic ESR intensity of the excited triplet states,  $I_{\text{int}}$ , which are the transitions  $|M = -1\rangle \rightarrow |M = 0\rangle$  and  $|M = 0\rangle \rightarrow |M = +1\rangle$ , is proportional to the radiation coefficient  $w$  and the difference of the population between the transition levels, i.e.

$$I_{\text{int}} = w \left\{ 2(N_{\text{int}, M=-1} - N_{\text{int}, M=0}) + 2(N_{\text{int}, M=0} - N_{\text{int}, M=+1}) \right\}. \quad (2)$$

The factor of 2 corresponds to the square of the transition matrix element.<sup>21</sup> When the intrinsic intensity  $I_{\text{int}}$  (i.e. intensity of  $\alpha$ ) is normalized by  $wN_{\text{imp}}$ , which was already obtained above, the unknown radiation coefficient  $w$  can be eliminated, and then we obtain

$$I'_{\text{int}} = \frac{N_{\text{int}}}{N_{\text{imp}}} \times \frac{2}{Z(T, B)} \times \sum_{\text{triplets } i} \left\{ e^{-\beta(E_i - g\mu_B B)} - e^{-\beta(E_i + g\mu_B B)} \right\}, \quad (3)$$

where  $Z(T, B)$  is the partition function

$$Z(T, B) = \sum_{\text{singlets } k} e^{-\beta E_k} + \sum_{\text{triplets } i} \left\{ e^{-\beta(E_i - g\mu_B B)} + e^{-\beta E_i} + e^{-\beta(E_i + g\mu_B B)} \right\}. \quad (4)$$

$N_{\text{int}}$  is the number of  $\{\text{Mo}_{75}\text{V}_{20}\}$ , and  $E_k$  and  $E_i$  are the energies of the singlet and triplet states, respectively. The normalized values of  $I'_{\text{int}}$  are presented as open circles in Fig. 3. The error bars are obtained from the uncertainty of the fitting curves. Using (3)  $I'_{\text{int}}$  can be evaluated exactly using the energy levels of the singlet and triplet states which can be obtained from the model Hamiltonian by diagonalization. This way the exchange constants can be determined by comparing experimental and theoretical values of  $I'_{\text{int}}$ .

It is important to note that the whole procedure works with only singlet and triplet levels since the exchange constants are large and therefore states with  $S \geq 2$  are irrelevant.

## B. Theoretical Models and analysis

The magnetism of  $\{\text{Mo}_{75}\text{V}_{20}\}$  is modeled by a Heisenberg Hamiltonian augmented with a Zeeman term

$$H = \sum_{i < j} J_{ij} \vec{s}_i \cdot \vec{s}_j + g\mu_B B \sum_j s_j^z. \quad (5)$$

Here  $J_{ij}$  is the exchange interaction between spins at sites  $i$  and  $j$ . The eigenvalues of this Hamiltonian can either be determined by complete matrix diagonalization or for the low-lying levels by the Lanczos method.<sup>22–24</sup> For our simulation we considered the coupling  $J_1$  on the equator and  $J_2$  to the tips of the triangles as schematically shown in the inset of Fig. 1. The low-energy part of the singlet and triplet sectors for  $J_2 = J_1$  (i.e. equilateral triangle) is shown in the inset of Fig. 3. Note that energies are given as multiples of  $J_1$ . The lowest singlet level is degenerate due to the frustration of the triangular magnetic structure.<sup>11</sup>

From the energy eigenvalues we can calculate  $I'_{\text{int}}$ . The ratio  $N_{\text{int}}/N_{\text{imp}} = 1/2$  is used since impurities of about  $2\mu_B$  per molecular unit are observed in the magnetization measurements. It is also important to note that we have taken into account all singlet and triplet levels up to  $10E_g$ , where  $E_g$  is the gap between the ground state and the lowest triplet level (see inset of Fig. 3). In view of the experimental temperature range this is more than sufficient since the exchange parameters turn out to be rather large.

The calculated  $I'_{\text{int}}$  curves for various  $J_1 = J_2$  are shown in Fig. 3. The position of the maximum of the calculated curves changes by varying the exchange coupling constant. The same holds true for the initial rise of the curve which is related to the energy gap between the singlet and triplet states, i.e.  $E_g$ . On the other hand, the high-temperature behavior of  $I'_{\text{int}}$  is in part related to the number of low-lying singlet states,<sup>25</sup> therefore, the high-temperature tail is almost independent of the exchange parameters. For  $J_1 = J_2 = 400$  K (thick solid curve in Fig. 3), the calculated curve fits well with the initial rising part and the tail part. However, a large difference is

seen on the peak position. This suggests that an equilateral triangle model ( $J_2 = J_1$ ) is not suited for this system, and that an isosceles triangle model ( $J_2 \neq J_1$ ) should be considered.

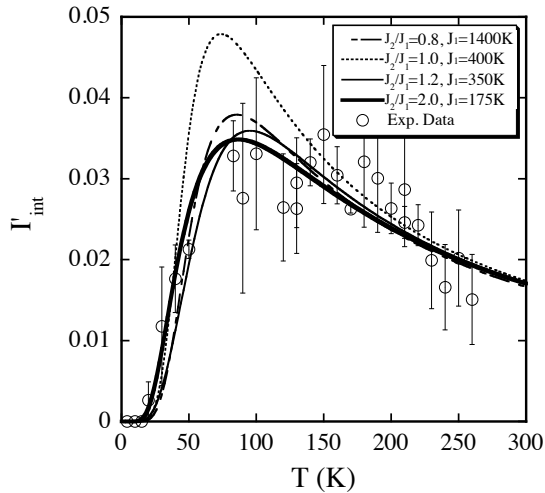


FIG. 4: Normalized integrated intensities for resonance  $\alpha$  coming from the intrinsic  $\{\text{Mo}_{75}\text{V}_{20}\}$  for various ratios  $J_2/J_1$ . The respective exchange parameters are shown in the legend.

Hence, we have investigated chains of isosceles triangles with various ratios  $J_2/J_1$ . Figure 4 shows the curves for  $J_2/J_1 = 0.8, 1.2, 2.0$ . The absolute values of the data are determined by the low-temperature behavior, i.e. the singlet-triplet gap. It is obvious that a much improved fit to our experimental data is provided by assuming the isosceles triangle model. The agreement in the observable  $I'_{\text{int}}$  appears to be better if  $J_2 > J_1$ . We emphasize that there is no need to introduce the additional exchange coupling  $J_3$  between tips of neighboring triangles. This exchange is very unlikely anyway due to the very long exchange pathway.

In order to rationalize the differences between  $J_1$  and  $J_2$  it is helpful to look at the crystal structure of an individual triangle of  $\{\text{Mo}_{75}\text{V}_{20}\}$  as schematically shown in Fig. 5. While the bonding length between vanadium atoms along  $J_1$  and  $J_2$  bonds does not vary much, the V-O-Mo bonding angles do. Moreover, since the V-O-Mo bonding angles are larger along  $J_2$  bonds the Goodenough-Kanamori rule<sup>26</sup> suggests that the exchange should be stronger along this pathway, which is in accord with our observations. Among the two good fits,  $J_2/J_1 = 1.2$ ,  $J_1 = 350$  K and  $J_2/J_1 = 2.0$ ,  $J_1 = 175$  K, we tend to favor the second parameter set since this reproduces the low-temperature behavior much better. The absolute numbers are in good agreement with other polyoxometalates containing  $\text{V}^{4+}$  spin centers.<sup>6,7,9</sup>

Our results suggest that the type of frustration which is present in  $\{\text{Mo}_{75}\text{V}_{20}\}$  is different from that in the original Keplerate molecules which are akin to the kagome lattice. Not only is  $\{\text{Mo}_{75}\text{V}_{20}\}$  a quasi one-dimensional object,

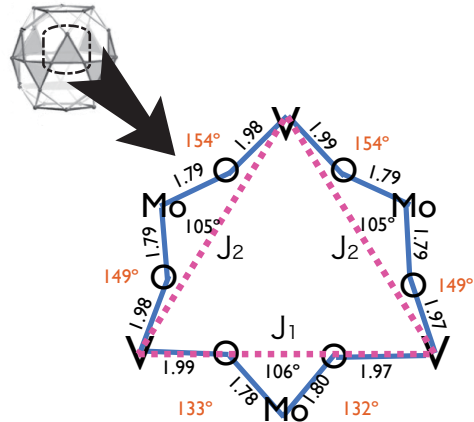


FIG. 5: Bond lengths and angles for an individual triangle of  $\{\text{Mo}_{75}\text{V}_{20}\}$ .

in addition the magnetic centers are not equivalent. The special ratio of  $J_2/J_1 = 2.0$  relates it strongly to sawtooth chains with flat bands of one-magnon energies.<sup>13,14</sup> Such chains are characterized by giant magnetization steps of 50 % of the saturation magnetization. Unfortunately, the exchange interactions present in  $\{\text{Mo}_{75}\text{V}_{20}\}$  are much too large in order to observe such magnetization steps experimentally, even for the first one, let alone subsequent steps at larger fields.

#### IV. SUMMARY

In summary, we have performed ESR measurements of the polyoxometalate cluster  $\{\text{Mo}_{75}\text{V}_{20}\}$ . We succeeded in separating the intrinsic ESR signal of the molecules and the signal of the magnetic impurities. Via the temperature dependence of the integrated intensity we could deduce the parameters of the underlying Heisenberg Hamiltonian. We also found that the additional exchange coupling  $J_3$ , which has been taken into account in previous studies<sup>3,17</sup>, is not necessary to explain the magnetic properties of  $\{\text{Mo}_{75}\text{V}_{20}\}$ .

Finally, we would like to stress that ESR is a powerful tool to study the low energy spectrum especially of frustrated magnetic systems. Since the sensitivity of ESR is very high, this method can in certain cases be far more effective than the specific heat measurements in characterizing the low-lying density of states. This study shows only ESR results for a single frequency of 190 GHz. It is of course possible to obtain more detailed information on the spectrum by combining results using multiple frequencies.

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